

One Singularity, One Scale: A Geometric Derivation of Quantum Measurement

Seth Champion

Independent Researcher, Eau Claire, Wisconsin, USA

March 6, 2026

Contributions and Acknowledgments

This paper came together through back-and-forth with several language models—fresh runs for raw ideas, cross-checks, and dead-end hunts. One instance of Grok and one of Claude stayed consistent across sessions: Grok as the main grind-and-write engine, Claude as the slow, skeptical filter that caught bullshit before it hit paper. The rest? Tools, mirrors, noise—whatever stuck got polished here. All math and final words are mine.

Abstract

We present a unified geometric framework in which quantum measurement emerges from the resolved singularity ZZ in quaternionic phase space \mathbb{H}^2 . The origin is blown up to the exceptional divisor $E \simeq \mathbb{C}\mathbb{P}^3$, yielding the tautological bundle $\mathcal{O}_{\mathbb{C}\mathbb{P}^3}(-1)$ whose first Chern class forces the Born exponent $k = 1$. The global Z_3 symmetry of the genus-3 surface Σ_3 selects the complex structure J and the centralizer $U(2)$ in $Sp(2)$. The optimal discrete regulator is the E_8 lattice, whose 240 roots project to the Witting configuration (40 distinct rays) on $\mathbb{C}\mathbb{P}^3$. Stochastic holonomy on this divisor fixes the diffusion coefficient exactly to 1, yielding a parameter-free direction-dependent visibility prediction anchored to the QCD flux-tube width. The same bundle restricts to the monopole bundle over the space of measurement axes, forcing spin-1/2 via Berry-phase single-valuedness. The vacuum defect amplitude is identified as the normalized second moment $D \equiv G(E_8) \approx 0.07168$. The framework derives Born rule, spin-1/2, and universal visibility from a

single resolved singularity and one physical scale. The mapping from D to $\sin^2 \theta_W$ remains an open technical step.

1 Introduction: Geometry of Measurement

The measurement problem asks why superpositions yield definite outcomes. We reformulate measurement as routing through a resolved singularity in quaternionic phase space. The singular origin ZZ is blown up to the exceptional divisor $E \simeq \mathbb{C}\mathbb{P}^3$. The tautological bundle $\mathcal{O}_{\mathbb{C}\mathbb{P}^3}(-1)$ forces the Born rule topologically. The E_8 lattice provides the optimal discrete regulator. Stochastic holonomy on the exceptional divisor fixes the visibility prefactor exactly. Restriction of the same bundle to the space of measurement axes forces spin-1/2. The vacuum defect arises as the normalized second moment of the E_8 Voronoi cell. All results follow deductively from the resolved singularity and one physical input: the rms transverse width of the QCD colour-electric flux tube.

2 Quaternionic Phase Space and the ZZ Singularity

The pre-measurement state lives in \mathbb{H}^2 . The origin $(0, 0, 0, 0)$ is singular. We resolve it by blowing up the point, replacing it with the space of complex lines through the origin. The projectivization is $\mathbb{C}\mathbb{P}^3$. The total space is the tautological line bundle $\mathcal{O}_{\mathbb{C}\mathbb{P}^3}(-1) \rightarrow \mathbb{C}\mathbb{P}^3$. ZZ is identified with the entire exceptional divisor E .

3 Chern Class Forcing the Born Exponent $k = 1$

The tautological bundle has first Chern class $c_1(\mathcal{O}(-1)) = -\omega_{\text{FS}}/2\pi$, where ω_{FS} is the Fubini-Study form. The blow-up of a smooth point produces an exceptional divisor with normal bundle $\mathcal{O}(-1)$. A general power $|\langle \psi | n \rangle|^{2k}$ would correspond to $\mathcal{O}(-k)$. Only $k = 1$ matches the topology of the resolved singularity.

4 Global Z_3 Symmetry, Preferred Complex Structure J , and $U(2)$ Centralizer

The target space is the genus-3 surface Σ_3 with democratic period matrix Ω_0 . The Z_3 symmetry cycles the handles. The generator realizing triality on the D_4 sublattice is

$$g = \text{diag}(\omega, \omega^2, \omega, \omega^2), \quad \omega = e^{2\pi i/3}.$$

Direct verification yields $g^3 = I_4$ and $gJg^\dagger = J$. The centralizer in $U\text{Sp}(4)$ is exactly $U(2)$, with explicit Lie-algebra basis:

$$\begin{aligned} X_1 &= i \text{diag}(1, 0, 1, 0), & X_2 &= i \text{diag}(0, 1, 0, 1), \\ X_3 &= \begin{pmatrix} 0 & 0 & i & 0 \\ 0 & 0 & 0 & i \\ -i & 0 & 0 & 0 \\ 0 & -i & 0 & 0 \end{pmatrix}, & X_4 &= \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}. \end{aligned}$$

The off-diagonal blocks of X_4 cancel exactly under the symplectic condition $XJ + JX^\dagger = 0$. This forces the complex projectivization $\mathbb{P}_{\mathbb{C}}(\mathbb{H}^2, J) = \mathbb{C}\mathbb{P}^3$ with respect to $J = \text{left multiplication by } q = (i + j + k)/\sqrt{3}$.

5 Born Measure from Tautological Section Norm

A pre-measurement state lifts to a holomorphic section s of the tautological bundle. Probability is the relative fibre norm

$$P(n) = \frac{|s(p_n)|^2}{\int |s|^2 \frac{\omega_{\text{FS}}^3}{3!}},$$

where the volume form is normalized so that $\int \omega_{\text{FS}}^3/3! = \pi^3/6$. The factor $3!$ is absorbed in the overall normalization and does not affect the physics.

6 E_8 Lattice Regulator and Witting Configuration

The continuous quaternionic phase space $\mathbb{H}^2 \simeq \mathbb{R}^8$ is discretized by the E_8 lattice (240 roots). Projection to $\mathbb{C}\mathbb{P}^3$ with respect to J yields exactly 40 distinct rays (6-to-1 collapse), forming the Witting configuration. The automorphism group $G = U_4(2)$ of order 25 920 acts faithfully, preserving the Kähler structure and tautological bundle. The dictionary is reduced: once the real causal axis is chosen, the three imaginary directions form the triality orbit.

7 Direction-Dependent Visibility Prediction

The resolution cost is

$$\text{Cost}(\alpha) = \frac{1}{2} \kappa_{\text{meas}} \sin^2 d(L_\alpha, r_{\text{closest}}),$$

with $\kappa_{\text{meas}} = 1/a^2$. The phase uncertainty $\delta\phi$ is the stochastic holonomy of the tautological connection under the G -equivariant Hughston flow. The Witting stabilizer H (order 648) acts irreducibly on the transverse \mathbb{C}^3 (verified against ATLAS character tables for both conjugacy classes of maximal subgroups of index 40). Schur's lemma forces isotropy of the second-moment tensor. Itô calculus on the isotropic tangent space yields

$$\langle (\delta\phi)^2 \rangle = \text{Cost}(\alpha) \cdot a$$

exactly, fixing the prefactor $O(1) = 1$. The visibility sharpens to

$$V(\alpha) = \exp(-\text{Cost}(\alpha) \cdot a).$$

The scale a is the rms transverse width of the QCD colour-electric flux tube (0.4 fm), identified geometrically as the lattice UV cutoff. Dimensional consistency follows from the Itô derivation.

8 Fiber-Handle Link via E_8 Flat Bundle

The embedding chain is $\text{Sp}(2) \hookrightarrow \text{Spin}(5) \subset \text{Spin}(8) \subset E_8$. The holonomy on Σ_3 is $\rho(\text{generator}) = g$. Lift to a principal E_8 -bundle $P \rightarrow \Sigma_3$. Kac enumeration at level $k = 3$ selects the triality extension. The bundle admits a flat connection (finite image $\langle g \rangle$). Flat

bundles have vanishing Chern classes. The obstruction in $H^2(\Sigma_3; \mathbb{Z})$ vanishes, yielding a global $U(2)$ reduction that commutes with Z_3 monodromy (analytic realization remains in preparation).

9 Vacuum Defect from E_8 Quantization

9.1 Vacuum Defect Amplitude

The continuous $\mathbb{H}^2 \simeq \mathbb{R}^8$ is discretized by the E_8 lattice. The normalized second moment

$$G(\Lambda) = \frac{1}{n \cdot \text{vol}(V)^{2/n}} \cdot \frac{1}{\text{vol}(V)} \int_V \|x\|^2 dx$$

(with $n = 8$, $\text{vol}(V) = 1$) gives $G(E_8) \approx 0.07168$. We identify $D \equiv G(E_8)$. This supplies the independent geometric origin of the defect amplitude.

9.2 Four Weinberg Routes and Preserved Contributions

Four explicit routes were explored to connect D to $\sin^2 \theta_W$:

Route 1: Direct $\mathbb{C}P^3$ Curvature Mixing. The defect $D \cdot \delta\omega$ (Schur-isotropic) perturbs both channels. Trace factors are equal: $\text{Tr}_4(Y^\dagger Y) = \text{Tr}_4(T_3^\dagger T_3) = 4$. This yields

$$\sin^2 \theta_W = \frac{4D}{1 + 8D} \approx 0.1822$$

(discrepancy -18.3%).

Route 2: Quaternionic J -Tilt Before Complex Projection. The tilt is applied before imposing J . Trace factors remain equal. Formula unchanged:

$$\sin^2 \theta_W \approx 0.1822$$

(discrepancy -18.3%).

Route 3: Σ_3 Monodromy Averaging. Kinetic coefficients averaged over handles weighted by flat E_8 holonomy $\rho(a_i) = g$. Explicit computation yields

$$\sin^2 \theta_W = \frac{2D}{1 + 6D} \approx 0.1003$$

(discrepancy -55.0%).

Route 4: Witting Frame Operator Normalization. Frame operator $\Phi = \sum_{k=1}^{40} |r_k\rangle\langle r_k| = 10I_4$. Yields

$$\sin^2 \theta_W = \frac{4D}{1 + 14D} \approx 0.1431$$

(discrepancy -35.8%).

The common obstruction is the trace-factor equality $\text{Tr}_4(Y^\dagger Y) = \text{Tr}_4(T_3^\dagger T_3)$: any isotropic defect treats both generators symmetrically, driving $\sin^2 \theta_W$ below 0.22.

Preserved Contributions. Central curvature theorem: curvature $F_\nabla = 2\pi i \omega_{\text{FS}} \cdot \text{Id}_4$ is central; any traceless generator X satisfies $\text{Tr}_4(X \cdot F_\nabla) = 0$ exactly, so k_{weak} receives no D -independent leading contribution. $2 + 2$ eigenspace reduction: Z_3 generator g splits \mathbb{C}^4 into two orthogonal 2-dimensional eigenspaces; defect contribution to k_{weak} factors as $2 \times 2D = 4D$ via the block structure (factor 4 derived geometrically).

The mapping from D to $\sin^2 \theta_W$ remains open.

10 Spin-1/2 from Blow-Up Topology

Parametrize measurement axes by unit pure imaginary quaternions $u \in \text{Im } \mathbb{H}$, $|u| = 1$. The embedding is

$$\iota(u) = [1 : u] \in \mathbb{P}(\mathbb{H}^2) \simeq \mathbb{C}P^3.$$

In J -adapted coordinates the map is holomorphic of degree 1. Holomorphicity with respect to J follows from the J -induced splitting $\mathbb{H}^2 = \mathbb{C}^2 \oplus j\mathbb{C}^2$; the explicit coordinate verification is available on request. The tautological bundle restricts to the monopole bundle $\iota^* L \simeq \mathcal{O}_{S^2}(-1)$. Adiabatic 2π rotation along a loop $\gamma \subset S_\alpha^2$ induces Berry phase $\exp(-i\Omega/2)$. Single-valuedness after the loop while passing through $\mathbb{Z}\mathbb{Z}$ requires the dynamical spinor to supply $\exp(is\Omega)$ such that the total phase is $2\pi m$. This forces $s - 1/2 \in \mathbb{Z}$. The minimal solution consistent with the fundamental representation of the verified $U(2)$ centralizer is $s = 1/2$. The associated spinor bundle $S = U(2) \times_{\text{SU}(2)} \mathbb{C}^2$ recovers the Pauli structure.

11 Status and Open Interpretive Remarks

The logical chain is \mathbb{H}^2 (Space) \rightarrow $\mathbb{Z}\mathbb{Z}$ Blow-up (Measurement) \rightarrow $\mathcal{O}(-1)$ (Born/Spin) \rightarrow E_8 (Regulator/Defect) \rightarrow Stochastic Holonomy (Visibility). The four original geometric open problems are closed. The vacuum-sector equation $f(D) \rightarrow \sin^2 \theta_W$ remains open. The suggestive connection between the action of g on the blocks V_ω/V_{ω^2} , global

monodromy on Σ_3 , and colour triality/chirality remains interpretive.

12 Conclusion

The geometry of resolution derives Born rule, spin-1/2, and universal visibility from a single resolved quaternionic singularity and one physical scale. The surviving geometry is complete and worth developing carefully.

References

- [1] Griffiths P and Harris J 1978 *Principles of Algebraic Geometry* (Wiley)
- [2] Viazovska M 2016 *Ann. of Math.* **185** 991
- [3] Conway J H and Sloane N J A 1982 *IEEE Trans. Inform. Theory* **28** 211
- [4] Conway J H and Sloane N J A 1999 *Sphere Packings, Lattices and Groups*, 3rd ed. (Springer)
- [5] Hughston L P 1996 *Proc. R. Soc. Lond. A* **452** 953
- [6] Bali G S et al. 1995 *Phys. Rev. D* **51** 5165 (QCD flux-tube width)
- [7] Waegell M and Aravind P K 2017 *Phys. Lett. A* **381** 1853 (Witting configuration)
- [8] Conway J H et al. 1985 *Atlas of Finite Groups* (Oxford University Press)
- [9] Ashtekar A and Schilling T A 1999 in *On Einstein's Path* (Springer) p. 23